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## THE CALCULATION OF THE HEAT CONTROL ACCUMULATOR VOLUME OF TWO-PHASE HEAT TRANSFER LOOP OF A SPACECRAFT THERMAL CONTROL SYSTEM

Spacecraft thermal control systems based on two-phase mechanically pumped loops have advantages in terms of mass and power consumption for auxiliary needs compared to single-phase thermal control systems. However, the disadvantage of two-phase mechanically pumped loops is that when changing the heat load and heat removal conditions, when switching from single-phase to two-phase operation mode and vice versa, the amount of working fluid in the loop changes significantly, which requires the use of a large volume heatcontrolled accumulator. Therefore, determining the minimum required volume of the heat-controlled accumulator for the loop operation is an urgent task due to the need to maintain the performance of the l loop at a minimum and maximum heat loads and minimize the mass of the structure and the working fluid charged. When determining the volume of the heat-controlled accumulator, it is necessary to correctly calculate the mass of the fluid in the loop during the two-phase operation mode. The mass of the fluid depends on the void fraction, which depends significantly on the phase slip. Many models and correlations have been proposed to calculate the phase slip factor. However, they all require justification for the parameters characteristic of spacecraft thermal control systems and weightlessness conditions. The paper presents the results of groundbased experiments, based on which the verification of different models and correlations for phase slip was performed. The validation of models and correlations for the conditions of weightlessness was performed by comparing the results with the horizontal and vertical orientation of the elements of the experimental setup. The working fluid is ammonia. The experiments showed that the best coincidence of calculation and experience is provided by Chisholm correlation. The discrepancy between the calculated and experimental values did not exceed +/-7% in the entire range of study parameters both for horizontal and vertical orientations, which allow us to recommend the Chisholm correlation for determining the coolant mass in the two-phase mechanically pumped loops for parameters characteristic of spacecraft thermal control systems, including zero-gravity conditions.

*Keywords:* spacecraft; weightlessness; thermal control system; two-phase mechanically pumped loop; heatcontrolled accumulator; volume; ammonia; slip factor.

### Introduction

For spacecraft (SC) with large heat dissipation (more than 6 kW), the most promising is thermal control systems (TCS) are based on a two-phase mechanically pumped loop (2PMPL) [1]. The NASA Technology Roadmap [2] classifies mechanically pumped two-phase loops as preferred for a series of future agency space missions. 2PMPL has many advantages compared to single-phase TCS, which, as a consequence, lead to significant weight savings of the system [3, 4]. A simplified schematic diagram of TPL with pumped fluid is shown in Fig. 1.

The loop consists of a pump (P), heat acquisition system (HAS), including thermal sinks/evaporators (EV), heat rejection system (HRS) with condensers (C), heat-controlled accumulator (HCA), liquid (Liquid), and vapor (Vapor) lines. The liquid work fluid is pumped by the pump (P) to the heat acquisition system (HAS), where it flows through the thermal sinks/evaporators (EV), where it is supplied with heat. As a result, the working fluid is heated to saturation, evaporates, and enters the two-phase state. The two-phase fluid enters the condensers (C) of the heat rejection system (HRS) where it is condensed and supercooled, as a result of heat removal to the environment, and then enters the pump through the liquid line (Liquid). The Heat Controlled Accumulator (HCA) is required to regulate the parameters of the loop as well as to compensate for changing the amount of fluid in the loop when the operating conditions of the loop change.

The most promising is the heat-controlled accumulator (HCA) [5-7].

As follows from the working principle, 2PMPL has a disadvantage compared with a single-phase loop: when changing the heat load and heat removal condi-

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tions, when switching from single-phase to two-phase mode and vice versa, the amount of working fluid in the loop changes significantly, especially in the case of a large volume of the two-phase section compared with the volume of the entire loop. Therefore, it is necessary to have a sufficiently accurate method of calculating the required volume of HCA, because if the volume of HCA is underestimated, the loop can lose control or fail at close to the limit operating modes, and overestimated volume of HCA will increase the weight of the system and the charged working fluid.



Fig. 1. Conceptual diagram of the 2PMPFL:
P – pump; HAS – heat acquisition system; EV – thermal sinks / evaporators; HRS- heat rejection system;
C – condensers; HCA – heat-controlled accumulator; Liquid – liquid line; Vapor – vapor line

### 1. Required HCA volume calculation

The required HCA volume  $V_{HCA}$  is equal to the volume of liquid working fluid that the accumulator must absorb or deliver to the loop, plus the guaranteed vapor pad volume in "cold" mode and the residual liquid volume in "hot" mode of operation of the loop. In this case, it is still possible to regulate the loop parameters using HCA, including the limit "cold" and "hot" modes. The mass of the fluid flowing in/out of the HCA is determined as the difference of the working fluid mass in the loop in single-phase mode  $M_{loop}^{cold}$  (minimum heat loads, "cold" orbit) and in two-phase mode  $M_{loop}^{hot}$  maximum heat loads, "hot" orbit).

The main calculation formula for the volume of HCA can be described as:

$$V_{\text{HCA}} = \left( M_{\text{loop}}^{\text{cold}} - M_{\text{loop}}^{\text{hot}} \right) / \left( \rho_{\text{m}}^{\text{hot}} - \rho_{\text{m}}^{\text{cold}} \right), \quad (1)$$

where:  $\rho_m^{hot}$  the average density of the two-phase mixtures in the HCA in the "hot" mode, kg/m<sup>3</sup>;  $\rho_m^{cold}$  the average density of the two-phase mixture in the HCA in the "cold" mode, kg/m<sup>3</sup>. Average density of the two-phase mixture in the HCA in the "hot" mode:

 $\rho_{m}^{hot} = \phi_{hot} \bullet \rho_{V} (t_{max}) + (1 - \phi_{hot}) \bullet \rho_{L} (t_{max}), \quad (2)$ where  $\phi_{hot}$  – ratio of vapor volume to full HCA volume in "hot" mode;

 $ho_V(t_{max})$  – vapor density at temperature  $t_{max}$  kg/m<sup>3</sup>;

$$\label{eq:rhoL} \begin{split} \rho_L(t_{max}) & - \text{ vapor density at temperature } t_{max} \\ kg/m^3; \end{split}$$

 $t_{\rm max}$  – temperature of the working fluid in the HCA at the maximum heat dissipation in the loop and "hot" orbit.

Average density of two-phase mixture in HCA in "cold" mode:

 $\rho_{m}^{cold} = \varphi_{cold} \bullet \rho_{V} (t_{min}) + (1 - \varphi_{cold}) \bullet \rho_{L} (t_{min}), (3)$ where  $\varphi_{cold}$  – ratio of vapor volume to full HCA volume in "cold" mode;

$$\label{eq:rho_V} \begin{split} \rho_V(t_{min}) & - \text{ vapor density at temperature } t_{min} \\ kg/m^3; \end{split}$$

 $\rho_L(t_{min})$  – vapor density at temperature  $t_{min}$  kg/m<sup>3</sup>;

 $t_{min}\,$  – temperature of the working fluid in the HCA at the maximum heat dissipation in the loop and "cold" orbit.

Formula (1) considers safety factors by setting  $\phi_{hot} > 0$  and  $\phi_{cold} < 1$  then calculating the average density of working fluid in the HCA.

The accuracy of the calculation of the required volume of HCA affects the performance of the circuit at the limit "cold" and "hot" modes and the possibility of regulating its parameters using HCA.

The main ways to reduce the required volume of HCA, based on the structure of formula (1), are as follows: it is necessary to increase  $M^{hot}_{loop}$  and  $\rho^{hot}_m$  decrease  $\,M^{cold}_{loop}\,$  and  $\,\rho^{cold}_{m}$  Analysis shows that the possibility of influencing the value  $M_{loop}^{cold}$  and  $M_{loop}^{hot}$  by changing the parameters of the working fluid (e.g., flow rate) for the loop of a given design at a given value of the maximum heat load and heat dissipation conditions, are limited. Therefore, the main way to reduce the HCA volume is to increase  $\,\rho_m^{hot}$  and the "hot" mode and decrease  $\rho_m^{cold}$  and the "cold" mode. In the limit it meats: to allow 100 % filling of HCA with liquid in the "hot" mode (  $\phi_{hot} = 0$  and complete emptying of HCA in the "cold" mode ( $\phi_{cold} = 1$ . But for this purpose, it is necessary to be able to calculate very accurately the necessary volume of heat-controlled accumulator.

At the stage of 2PMPL design, there are uncertainties and errors in calculating the mass of the fluid in the loop. Especially great is the error in calculating the mass  $M_{loop}^{hot}$  the two-phase mode at the maximum heat load and the worst heat removal conditions. This forces the entry of increased reserve coefficients by setting  $\phi_{hot} > 0$  and  $\phi_{cold} < 1$  which negatively affects the volume and total mass of the heat-controlled accumulator and TCS.

# 2. Working fluid mass calculation in the loop for two-phase mode

The mass of the working fluid in the loop is equal to the sum of the masses in the individual sections. To calculate the total mass of the working fluid, the entire 2PMPL must be divided into elementary sections, within which the flow parameters can be assumed to be unchanged.

Calculation of the working fluid mass in sections with single-phase liquid working fluid (liquid lines, pump, etc.) is not difficult, it is enough to know the temperature of the working fluid in these elements and the geometric parameters of the loop [8, 9].

Calculation of working fluid mass in sections with two-phase working fluid (evaporators, vapor lines, and condensers) presents significant difficulties since the mass is determined not only by the geometry and temperature of the working fluid but also by the distribution of mass and volume vapor quality, the slip of phases along the length of the channel.

The mass of the two-phase working fluid in an elementary section of a cylindrical pipeline of length  $\Delta X$  can be estimated by the formula:

$$\Delta M_{2p} = \Delta V_{2p} \bullet \rho_m , \qquad (4)$$

where  $\Delta V_{2p} = \mathbf{F} \cdot \Delta \mathbf{X}$  - volume of the elementary section

of the loop with two-phase working fluid, m<sup>3</sup>;

F- cross-sectional area of the channel.  $m^2$ ;

 $\Delta X$  – length of elementary section, m;

 $\rho_m$  – average density of two-phase working fluid in this section, kg/m<sup>3</sup>.

The average density of the two-phase working fluid depends on the phase densities and the void fraction.

$$\rho_{\rm m} = \rho_{\rm L} - \alpha \cdot (\rho_{\rm L} - \rho_{\rm V}), \qquad (5)$$

where  $\alpha = 1/(1 + \frac{1-x}{x} \cdot \frac{\rho_V}{\rho_L} \cdot S)$  the void fraction.

In the equations:

 $\rho_L$  and  $\rho_V$  – density of liquid and vapor phases, kg/m<sup>3</sup>,

x – vapor quality;

$$S = \frac{u_V}{u_L} - \text{slip factor;}$$

 $u_V$  and  $u_L$  – average velocity of liquid and vapor phases, m/s.

Thus, in order to correctly estimate the mass of the working fluid in the loop, along with other parameters, if is necessary to know the void fraction, which, in turn, depends significantly on the slip factor S.

Various models have been proposed to estimate the slip factor.

- Homogeneous flow for which the slip factor [10]:

- Momentum Flux Model [11]:

$$\mathbf{S} = \left(\frac{\rho_{\rm L}}{\rho_{\rm V}}\right)^{1/2};\tag{6}$$

- Zivi Kinetic Energy Model for annular flow [12]:

$$\mathbf{S} = \left(\frac{\rho_{\mathrm{L}}}{\rho_{\mathrm{V}}}\right)^{1/3};\tag{7}$$

 Levy Momentum Model [13]. The void fraction is defined by the equation solving:

$$x = \frac{\alpha (1 - 2\alpha) \alpha \sqrt{(1 - 2\alpha)^2 + \alpha (2 \frac{\rho_L}{\rho_V} (1 - 2\alpha)^2 + \alpha (1 - 2\alpha)}}{2 \frac{\rho_L}{\rho_V} (1 - 2\alpha)^2 + \alpha (1 - 2\alpha)}; (8)$$

- Chisholm Method [14]:

$$S = (1 - x(1 - \frac{\rho_L}{\rho_V}))^{1/2}; \qquad (9)$$

- Model proposed by M. Narcy et al. [15], in according with this model:

$$\alpha = \frac{\mathbf{f} \cdot \mathbf{x}^n}{1 + (\mathbf{f} - 1)\mathbf{x}^n},\tag{10}$$

where coefficients  $f = a + (1-a)(\frac{\rho_V}{\rho_V})^{a_1}$  and

$$n = b + (1-b)(\frac{\rho_V}{\rho_L})^{a_1}$$
 with values:  $a = -2.129$ ,

b = 0.3487,  $a_1 = -02186$ ,  $b_1 = 0.515$ .

As the authors point out, this model describes well the void fraction in microgravity conditions.

As an example, lets calculate the mass of ammonia in a section of adiabatic pipeline with a diameter of 7 mm and a length of 13.7 m. The flow parameters are as follows: vapor quality - 0.5; saturation temperature -65°C. Calculation result: if there is no slip (S=1)  $\Delta M_{2p} = 0.047$  kg; calculating the slip factor by Chisholm's model  $\Delta M_{2p} = 0.097$  kg.

The conclusion is clear: the slip of the phases must be included when calculating the mass of the working fluid in the loop in the two-phase modes of its operation.

# 3. Model verification for calculating the slip factor

The slip factor S depends on many mode parameters and channel geometry. The applicability of slip modeling proposed by various authors to calculate slip under specific conditions requires, as a rule, experimental verification.

Therefore, a set of experiments was performed to determine an acceptable model for calculating the slip of phases in an experimental two-phase heat transfer loop at different orientations of the evaporator branch. At the same time, the same slip model was used to calculate the working fluid mass in all two-phase sections of the loop.

The mass difference in the two-loop operating modes (single-phase and two-phase) was also calculated theoretically. The discrepancy between the calculated and experimental determination is the main criterion for the applicability of one or another calculation model for determining the mass of the working fluid in 2PMPL.

### 3.1. Experimental setup description

Fig. 2 shows a diagram of the experimental loop, on which the study was conducted. The working fluid is ammonia.



Fig. 2. Conceptual diagram of the experimental loop

The experimental setup was a single-loop loop containing pump (Pump), heater (HDH), single-phase transport sections, evaporator (Evaporator), two-phase transport section, condenser (Condenser). The pressure in the loop is regulated by a heat-controlled accumulator (HCA). All elements are located in the same horizontal plane. The evaporator branch could take horizontal or vertical position.

The evaporator is a heat exchanger with 10 heat sinks connected in series with a channel diameter of 7 mm, on which 10 heaters with a total capacity of 6 kW are installed. The total length of the evaporator is 13700 mm.

The two-phase pipeline connects the evaporator and condenser and is designed as a tube with a diameter of 7 mm and a length of 7000 mm.

The condenser is a connected two-channel aluminum profile. One channel is filled with two-phase ammonia and the other with coolant glycol. The condensation section is significantly shorter than the other twophase sections.

The HCA was mounted on a high-precision weight scale.

### **3.2. Experimental procedure**

At first, a series of experiments were performed in the two-phase mode. The loop was set to stationary mode with the following specified parameters:

- Saturation temperature in HCA T<sub>HCA</sub>;
- mass flow rate of working fluid m;
- subcooling at the evaporator branch inlet  $\Delta T_{sub}$ ;
- glycol temperature T<sub>glic</sub>;
- total supplied power Qel.

The main part of the experiments was performed with all evaporator heaters turned on simultaneously.

After reaching the steady-state mode, all parameters of the 2PMPL were recorded, including the HCA mass in the two-phase mode  $M_{HCA_2f}$  and the total injected electric power Qel.

After conducting a series of experiments in twophase mode, all evaporator heaters (except the heater (HDH)) were switched off, the 2PMPL was switched to single-phase mode, while the flow rate and temperatures along the length of the single-phase part of the loop were maintained. All parameters of the 2PMPL were fixed, including the mass of the HCA in single-phase mode  $M_{HCA_{1f}}$ . Thus, the mass of fluid in "single-phase" elements of 2PMPL (pump, transport liquid pipeline) in "two-phase" and "single-phase" modes of 2PMPL operation remained unchanged.

Mass difference:

$$\Delta M_{\text{HCA}} = M_{\text{HCA}_2f} - M_{\text{HCA}_1f} \qquad (11)$$

determines the change of fluid mass in the loop only in "two-phase" elements of 2PMPL  $\Delta M_{MPL}$ : in evaporator branch, transport two-phase pipeline and condenser:

$$\Delta M_{\text{HCA}} = M_{\text{MPL}}.$$
 (12)

In other "single-phase" 2PMPL elements, the change in mass is negligible, since the temperature of the liquid fluid remained approximately constant.

The error in determining the heat supplied to the evaporators was determined in special calibration experiments and was as follows:

- at Tsat = 55 °C - no more than 30W;

- at Tsat = 75 °C - no more than 40W.

The main error in calculating the vapor quality of  $x_{ex}$  behind the evaporator is the amount of heat supplied to the fluid. Since the flow rate and electrical power were measured with high accuracy, the error is mainly determined by the error in calculating heat losses.

- for flow rates greater than 2 g/s, the relative error in determining the vapor quality was: 4...5 % for x $\approx$ 0,4; 2...3 % for x $\approx$ 0,8.

– for flow rates less than 2 g/s the relative error of vapor quality determination was 5...14 %.

The mass difference of the fluid in the HCA  $\Delta M_{HCA}$  is determined in experiments with high accuracy. The division value of the used scales is 1 g. The relative error of determining the fluid mass in the two-phase part of the loop did not exceed 1.8 %.

# **3.3.** Working fluid mass calculation in two-phase part of the loop

To calculate the mass of working fluid in the twophase part of the loop, a mathematical model of stationary flow distribution in the heat transfer loop was used. The ammonia mass in the two-phase section of the loop depends on the temperature, vapor quality distribution along the length of the channel, and slip between the phases. The heat input configuration in the evaporator was taken into account. The models for determining the slip factor of the phases were set the same in all twophase elements of the loop.

# 3.4. Experimental results and comparison with calculations

Six series of experiments on ammonia were performed. The range of parameter changes in the experiments was as follow:

– saturation temperature in HCA  $T_{HCA} = 55^{\circ}C$  and 75 °C;

- mass flow rate m = 0,8 ... 16 gm/s;

- vapor quality at evaporator branch outlet  $x_{ex} = 0,12 \dots 0,8;$ 

– subcooling at the evaporator inlet  $\Delta T_{sub} \sim 5...25$  K;

- glycol temperature  $T_{glic} \sim 20...50$  °C.

During changes in the heat load, the main change in the mass of the working fluid occurs in the elements with two-phase loop: in the evaporator, the transport section and in the condenser. Therefore, in the following we consider only the two-phase sections of the loop. The volumes of the "two-phase" elements of the experimental 2PMPL were as follows:

- total volume of two-phase part of 2PMPL is 1,207 L (100%);

- evaporator volume is 0,556 L (46,1%);

volume of two-phase transport section is 0,57 L (47,2%);

– volume of the condensation section ~ 0,081 L (6,7%).

Comparison of calculated and experimental values of coolant mass changes in the two-phase section of the loop during "single-phase" and "two-phase" modes served as a criterion for the applicability of this or that correlation to solve the problem of calculating the ammonia mass in the 2PMPL. In Fig. 3 a comparison was made when using the Chisholm model in the mathematical model of 2PMPL and different orientations of the evaporator.

A comparison of calculated and experimental values using other models shows that the Chisholm model provides the best match. The mismatch between the calculated and experimental values of the change in mass does not exceed +/-7% in the entire range of parameters of the experiments conducted for both horizontal and vertical orientations. The Chisholm model can also be recommended for calculating the working fluid mass in 2PMPL under weightlessness conditions since it has shown the best agreement between the calculated and experimental results in ground experiments in a wide range of parameters characteristic of TCS and for various evaporator orientations.

To select the best phase slip model, the parameter was also used:

$$\mathbf{RMS} = \sqrt{\frac{1}{N} \sum_{i=1}^{N} \left( \frac{\Delta M_{\text{exp},i} - \Delta M_{\text{sim},i}}{\Delta M_{\text{exp},i}} \right)^2} , \qquad (13)$$

where N - the number of experimental points;

 $\Delta M_{exp}$  – mass change of the working fluid in the HCA in the experiment;

 $\Delta M_{sim}$  – estimated mass change of the working fluid in the loop.

The RMS value for different series of experiments, saturation temperatures, evaporator orientation, and the model used to calculate the slip is shown in Table 1.

A comparison of the calculated and experimental values of the change in masses according to the RMS parameter shows that:

 Chisholm model provides the best match between the calculated and experimental values of the mass change for the horizontal orientation;

- for the vertical orientation, the best match is given by the Zivi and Levy models. However, for the verti-

Table 1

cal orientation the RMS discrepancy for the Zivi, Levy, and Chisholm models is insignificant. Therefore, it is recommended to use the Chisholm model for both horizontal and vertical orientations of the evaporator for convenient processing of the results of ground-based experiments.



Fig. 3. Comparison of calculated and experimental changes in mass in the loop in single-phase and two-phase modes. Chisholm correlation is used

RMS	value	for	different	models

	Horiz	ontal	Vertical	
T sat, °C	55	75	55	75
Homogeneous	0.135	0.163	0.197	0.219
Momentum flux	0.156	0.153	0.088	0.076
Zivi	0.062	0.054	0.049	0.077
Levy	0.065	0.049	0.043	0.073
Chisholm	0.029	0.039	0.052	0.09
Cincinatti	0.073	0.112	0.1	0.168

#### Conclusions

To maintain the performance of the two-phase mechanically pumped loop of the spacecraft thermal control system at a minimum and maximum heat load, to ensure the possibility of controlling the pressure in the loop over the entire operating range using a heatcontrolled accumulator, to minimize the weight of the structure and volume of the heat-controlled accumulator, it is necessary to calculate its required volume correctly.

A method of calculating the required volume of the heat-controlled accumulator is proposed. The calculation accuracy of the accumulator volume is determined mainly by the calculation error of the working fluid mass in the two-phase part of the loop during its operation in the "hot" mode (maximum heat load, the worst heat removal conditions). In addition to considering the real configuration of heat supply and heat removal and vapor quality distribution along the loop, the phase slip factor, which depends on the channel geometry and flow parameters, should be considered in the calculation model. Various models for calculating the phase slip factor have been analyzed (Homogeneous model, Momentum flux model, Zivi model, Levy model, Chisholm model, Cincinnati model). Calculation of fluid mass difference in the loop during single-phase "cold" mode (minimum heat loads, best heat removal conditions) and two-phase "hot" mode was compared with the experiment on the loop model. It was shown that the best coincidence of calculation and experiment takes place when the Chisholm model is used in calculations.

The single-loop experimental setup included all main elements of the standard system: pump, evaporator, two-phase transport section, condenser, and heatcontrolled accumulator. The experiments were carried out with natural working fluid (ammonia), at parameters characteristic of two-phase mechanically pumped loops of unmanned spacecraft thermal mode systems, at relatively high flow rates (when gravity does not affect hydrodynamics and heat exchange), at the horizontal and vertical location of the evaporator. This allows us to recommend using the Chisholm model to calculate phase slip in weightless conditions as well.

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### РОЗРАХУНОК ОБСЯГУ ГІДРОАКУМУЛЯТОРА ДВОФАЗНОГО КОНТУРУ ТЕПЛОПЕРЕНОСУ СИСТЕМИ ТЕРМОРЕГУЛЮВАННЯ КОСМІЧНОГО АПАРАТУ

### А. М. Годунов, Г. О. Горбенко, П. Г. Гакал

Системи терморегулювання космічних апаратів на базі двофазних контурів теплопереносу мають переваги по масі і енергоспоживанню на власні потреби в порівнянні з однофазними системами терморегулювання. Однак, недоліком двофазних контурів теплопереносу є те, що при зміні теплового навантаження і умов тепловідведення, при переході від однофазного до двофазного режиму роботи і навпаки, кількість теплоносія в контурі істотно змінюється, що вимагає використання гідроакумулятора великого обсягу. Тому, визначення мінімально необхідного для роботи контуру обсягу гідроакумулятора є актуальним завданням з огляду на необхідність збереження працездатності контуру при мінімальних і максимальних теплових навантаженнях і мінімізації маси конструкції і заправленого теплоносія. При визначенні обсягу гідроакумулятора необхідно коректно розраховувати масу теплоносія в контурі теплопереносу на двофазному режимі роботи. Маса теплоносія залежить від об'ємного паровмісту, який істотно залежить від ковзання фаз. Запропоновано багато моделей і кореляцій для розрахунку ковзання фаз. Однак, всі вони вимагають обгрунтування для параметрів, характерних для систем терморегулювання космічних апаратів і умов невагомості. У статті представлені результати наземних експериментів, на основі яких проводилася верифікація різних моделей і кореляцій для ковзання фаз. Обгрунтування застосовності моделей і кореляцій для умов невагомості проводилася шляхом порівняння результатів при горизонтальній і вертикальній орієнтації елементів експериментальної установки. Теплоносій-аміак. Експерименти показали, що найкращий збіг розрахунку і досвіду забезпечує кореляція Chisholm. Розбіжність між розрахунковим і експериментальним значеннями не перевищувала +/-7% у всьому діапазоні досліджених параметрів як для горизонтальної, так і для вертикальної орієнтації, що дозволяє рекомендувати кореляцію Chisholm для визначення маси теплоносія в двофазному контурі теплопереносу для параметрів, характерних для систем терморегулювання космічних апаратів, в тому числі і для умов невагомості.

Ключові слова: космічний апарат; невагомість; система терморегулювання; двофазний контур теплопереносу; гідроакумулятор з тепловим регулюванням; об'єм; аміак; ковзання фаз.

### РАСЧЕТ ОБЪЕМА ГИДРОАККУМУЛЯТОРА ДВУХФАЗНОГО КОНТУРА ТЕПЛОПЕРЕНОСА СИСТЕМЫ ТЕРМОРЕГУЛИРОВАНИЯ КОСМИЧЕСКОГО АППАРАТА

#### А. М. Годунов, Г. А. Горбенко, П. Г. Гакал

Системы терморегулирования космических аппаратов на базе двухфазных контуров теплопереноса обладают преимуществами по массе и энергопотреблению на собственные нужды по сравнению с однофазными системами терморегулирования. Однако, недостатком двухфазных контуров теплопереноса является то, что при изменении тепловой нагрузки и условий теплоотвода, при переходе от однофазного к двухфазному режиму работы и наоборот, количество теплоносителя в контуре существенно изменяется, что требует использования гидроаккумулятора большого объема. Поэтому, определение минимально необходимого для работы контура объёма гидроаккумулятора является актуальной задачей ввиду необходимости сохранения работоспособности контура при минимальных и максимальных тепловых нагрузках и минимизации массы конструкции и заправленного теплоносителя. При определении объема гидроаккумулятора необходимо корректно рассчитывать массу теплоносителя в контуре теплопереноса на двухфазном режиме работы. Масса теплоносителя зависит от объемного паросодержания, которое существенно зависит от скольжения фаз. Предложено много моделей и корреляций для расчета скольжения фаз. Однако, все они требуют обоснования для параметров, характерных для систем терморегулирования космических аппаратов и условий невесомости. В статье представлены результаты наземных экспериментов, на основе которых проводилась верификация различных моделей и корреляций для скольжения фаз. Обоснование применимости моделей и корреляций для условий невесомости проводилась путем сравнения результатов при горизонтальной и вертикальной ориентации элементов экспериментальной установки. Теплоноситель – аммиак. Эксперименты показали, что наилучшее совпадение расчета и опыта обеспечивает корреляция Chisholm. Расхождение между расчетным и экспериментальным значениями не превышало +/-7% во всем диапазоне исследованных параметров как для горизонтальной, так и для вертикальной ориентации, что позволяет рекомендовать корреляцию Chisholm для определения массы теплоносителя в двухфазном контуре теплопереноса для параметров, характерных для систем терморегулирования космических аппаратов, в том числе и для условий невесомости.

Ключевые слова: космический аппарат; невесомость; система терморегулирования; двухфазный контур теплопереноса; гидроаккумулятор с тепловым регулированием; объем; аммиак; скольжение фаз.

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